Temperature and Power Dependence of Photoluminescence in PbS Quantum Dots Nanoparticles

(Mesaranan Suhu dan Kuasa Penguajaan terhadap Fotoluminesens Titik Kuantum PbS Berzarah Nano)

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ABSTRACT
In this study, the synthesis and the effect of temperature and power excitation towards photoluminescence (PL) emission of colloidal PbS quantum dots (QDs) were reported. Water soluble PbS QDs capped with a mixture of 1-thioglycerol (TGL) and dithioglycerol (DTG) was synthesized via colloidal chemistry method at room temperature. The PL emission of PbS QDs was investigated under temperature range from 10 K to 300 K and we found that the PL emission blue-shifted when the temperature is increased. From high resolution transmission electron microscopy (HRTEM), the average size of PbS core QDs is determined to be 6 nm and the integrated PL intensity (IPL) versus excitation power density shows the recombination of electrons and holes occur efficiently at low and high temperature for the PbS QDs. Full width half maximum (FWHM) shows a gradual broadening with the increasing temperature due to the interaction of charge carriers with phonons.

Keywords: Near infrared; PbS; photoluminescence; quantum dots

INTRODUCTION
Colloidal quantum dots (QDs) have become one of the primary attention highlighted in current researched topic due to its unique intrinsic properties and potential used in optoelectronics and biomedical applications. Generally, QDs are semiconductor crystals with a size in the nanoscale range. It possesses a physical dimension that smaller than the exciton Bohr radius which lead to its excellent electronic properties based on the quantum confinement effect (Gao 2011). In this study, lead sulphide (PbS) QDs are selected because of its narrow bulk energy band gap at 0.41 eV and large exciton Bohr radius of 18 nm as compared to other materials such as CdS (2.8 nm) (Shamsudin & Junas 2018). It is worth mentioning that the effective mass of electron and holes in QDs are almost similar with $m_e^* = m_h^* \approx 0.08 m_0$ whereby $m_0$ is the electron mass in vacuum. This results eventually contributed to a strong confinement of both electrons and holes in the PbS QDs and benefit it potential in various optoelectronic application such as near infrared (NIR) photodetector, solar cell device, and light emitting diode (LED) (Guchhait 2011; Iacovo 2016; Reilly 2014; Zhou et al. 2019). In addition, the size dependent optical absorption coefficient of QDs also allowed it to be used as a switchable device and ease the spectra modification in absorption, photoluminescence and electroluminescence spectra (Okamoto 2004).

Various methods have been conducted to synthesis a QDs such as molecular beam epitaxy (MBE), metal organic chemical vapour deposition (MOCVD), hydrothermal method and colloidal chemistry (Ashari et al. 2016; Tasco 2007; Yang et al. 2006; Yoffe 2001). It is noteworthy that a high crystalline QDs can be fabricated by MBE and MOCVD, however, it requires a high vacuum-technology equipment and high cost process. In contrast, the wet chemical colloidal approach provides a simple synthesis, feasible and low cost process (Yin & Alivisatos 2005). It gives advantage to QDs because there is no need for QDs to be attached to a surface as they are prepared in solution. This can be used as nanometer size building blocks for interesting structures and electronic devices (Sargent 2005). Basically, the organic ligands were used during the colloidal technique, aiming to prevents the
agglomeration of PbS QDs due to the high surface free energy. Organic ligands also play a vital role in controlling the rate of kinetic growth of semiconductor crystal in a colloidal solution and indirectly stabilize the shape and dimension of QDs (Binetti et al. 2015; Rogach et al. 2007). Thus, the optimum selection of organic ligand is essential for surface functionalization, solubilization of ions during the preparation and modification of carrier confinements in QDs (Alivisatos 1996).

To understand the synergy relation between the localized states and carrier transport occurred in a PbS QDs based semiconductor material, temperature dependent photoluminescence spectroscopy is a common equipment used in the research. A study related to power excitation and temperature are introduced. There are many physical parameters that can be detected such as defects related to localized states, non-radiative processes, and carrier transport dynamics (Okamoto et al. 2004; Ryu et al. 2009; Szendrei et al. 2012; Yamada et al. 2009). Interestingly, the PbS QDs synthesized in this study present a high PL linewidth in the range of 150 meV to 170 meV when the operating temperature is increases from 10 K to 300 K. This results are outstanding as compared to previous report which detected at ~200 meV for PbS QDs embedded in glass (Peterson & Krauss 2006). It is strongly believed that the increases of excitation power density can bring forward the modification the electronic and optical properties in QDs and acted as a new approach to study the mechanism behind the exciton-exciton interactions.

In this report, we report a study of the electronic and optical properties of PbS QDs prepared via colloidal synthesis. The characteristic of the emission spectra was evaluated in the function of temperature and power excitation from the temperature range of 10 K to 300 K. The PL study showed that the samples possess a near infrared (IR) emission at the range of 1000 nm to 1300 nm. The dependency between the full width at half maximum (FWHM) and integrated PL intensity (IPL) of the sample are also demonstrated.

EXPERIMENTAL SECTION

MATERIALS

1-thioglycerol (TGL, 95%), dithioglycerol (DTG, 95%) were obtained from Sigma Aldrich, lead (II) acetate trihydrate (99.99%), sodium sulphide nonahydrate (99.99%), were obtained from R&M Chemicals, triethylamine was purchased from Chemiz. The deionized water was used throughout the experiments with resistance of 18.2 MΩcm⁻¹. All chemicals were used as purchased.

SYNTHESIS OF COLLOIDAL PbS QDS

The synthesis of PbS QDs in aqueous solution was stabilised with capping ligands, namely, TGL and DTG. The molar ratio of Pb²⁺ and S²⁻ precursors is 1 and 0.3, respectively. During the synthesis of PbS QDs, lead acetate salt in solid crystal form was dissolved in deionized water and oxygen was removed from the system by flowing nitrogen gas. This oxygen-free environment avoids the oxidation of the QDs and sedimentation phenomena (Wang et al. 2015). Meanwhile, 0.01 M solution of lead (II) acetate was prepared in a three-neck flask. 1 mL of the lead solution was added to the precursor beaker and the solution was allowed to mix for a few minutes. TGL and DTG were added into the three-neck flask containing lead precursor and magnetic stirrer running at 400 rpm. The solution then turned to yellowish green. These ligands (TGL and DTG) were added to avoid agglomerations. Triethylamine was added to the solution drop wise to adjust the pH until it reaches to pH10 and the solution turned colourless. After 15 min, 0.1 M solution of sodium sulphide was immediately introduced into the system, accompanied by continuous stirring. The solution changed from colourless to dark brown immediately, indicating the formation of lead sulphide QDs. All the samples were stored in the dark at low temperature (4°C).

CHARACTERIZATION

For optical measurements (PL spectroscopy), the 532 nm laser beam was used as an excitation source with the closed cycle He cryostat system measured temperature range from 10 K to 300 K as shown in Figure 1. The temperature controller was used to vary the temperature. The emitted PL was collected by a lens and then focused onto the entrance slit of the monochromator before being detected by an InGaAs detector. The integration time was 1 s and the investigated wavelength range from 800-1600 nm. For the sample preparation, the colloidal PbS QDs was drop casted on the glass substrate for PL measurement.

For high resolution field emission transmission electron microscope (HRTEM), the sample was prepared by drop-casting the small volume of the PbS QDs onto the Cu grid coated with an ultrathin coated film. The structural characterization was performed using FEI Tecnai-G2 20 S-TWIN TEM by Crest Group. The size distribution, pattern and particle size of QDs for the sample was characterized by HRTEM.

RESULTS AND DISCUSSION

COLLOIDAL SYNTHESIS

Figure 2 displays the colloidal PbS QDs synthesized in aqueous solution and capped with thiols ligands (TGL and DTG). These organic ligands have functional group of sulphur and hydrogen atom, which bind to the QDs surface with hydrogen bond. From Figure 2, the diagram shows the capping ligands preventing the PbS QDs from agglomeration by creating a protection to QDs from binding effect, where the area between QDs is covered by ligands, causing repulsion between QDs that stabilizes the colloid (Liu et al. 2018). Generally, long chain thiols are excellent capping ligand for most semiconducting and metal
nanoparticles. For instance, thiol stabilised colloidal CdTe QDs was firstly reported that led to the repetitive use of the ligands with QDs (Rajh et al. 1993). Unfortunately, this colloidal PbS QDs is highly sensitive to the environment and temperature. It was observed that the sample was agglomerated when exposed to ambient oxygen at relatively high temperature. Therefore, the colloidal PbS QDs was kept at low temperature to slow down the growing process of QDs, thus making the size of QDs smaller as well as preventing the formation of bulk crystals or molecular clusters instead of the anticipated QDs. It is also important to store the samples at relatively low temperature because at low temperature, it provides strong protection to the core of PbS QDs against the Ostwald ripening which allows us to attain uniform size distribution (Moreels et al. 2009).

STRUCTURAL PROPERTIES
Figure 3 shows the representative HRTEM image of the colloidal PbS QDs prepared in aqueous solution. From the Figure 3, it can be clearly observed that the PbS QDs are closely spherical in shape with monodispersity and clear lattice structure inside the QDs. The QDs size is approximately 6 ± 1 nm which is smaller than the exciton Bohr radii (18 nm), leading to strong quantum confinement effect by confining the motion of electron and hole. In addition, the QDs are clearly separated from binding together due to protection from the thiols ligands that provide separation distance between QDs. The thiols have high affinity for Pb²⁺ ions and protect the PbS QDs with a steric shielding that screens the Van Der Waals attractions between nearby QDs (Yu et al. 2013).

PHOTOLUMINESCENCE
TEMPERATURE DEPENDENCE OF THE COLLOIDAL PbS QDS PL EMISSION
Figure 4 shows the PL spectra for PbS QDs at a range of T (10 K–300 K). The increment of temperature has several effects on the PL peak energy, PL linewidth and the peak intensity. Based on Figure 4(b), PL peak energy was red-shifted when the temperature is increased from 20 K to 60 K, while it blue-shifted from T >80 K. The energy shifts for temperature higher than 80 K can be described by a temperature coefficient, $\alpha = \frac{dE}{dT} = 0.1$ meV/K. This
value is smaller compared to bulk PbS ($\alpha = 0.52$ meV/K), but is comparable to that reported in previous study of thiol-capped PbS QDs ($\alpha = 0.3$ meV/K) and PbS QDs in glass ($\alpha = 0.2$ meV/K) (Litvyak et al. 2017; Turyanska et al. 2007). Usually, the temperature dependence of the PL peak position can well be described by Varshni equation (Varshni 1967),

$$E_g(T) = E_g(0) + \frac{\alpha T^2}{\beta + T}$$

(1)

where $(0)$ is the band gap at $T = 0$ K; $\alpha$ represents the band gap energy temperature coefficient; and $\beta$ is a constant of the order magnitude of the semiconductor material Debye temperature. For this colloidal PbS QDs sample, the second term in (1) with the plus sign (rather than minus sign, as it was in the original equation) was used because the PbS QDs band gap energy increases with the increasing temperature, which is in contrast to most of semiconductor materials.

Commonly, the energy shifts of PL emission with temperature can be caused by some mechanisms. These include thermal expansion of the crystal lattice and electron-phonon coupling (Awschalom & Johnston-Halperin 2001; Wise 2000). When the temperature increases, the distance between atoms in QDs also increases. This will make the energy band gap rises with the increasing temperature as seen in the Figure 4(b) at $T > 80$ K of our PbS QDs sample. Another dominant contribution to the temperature variation of the energy band gap of the QDs is due to strain between the thermal expansion mismatch and between the QDs and its capping layer (Klotz et al. 2010). These complex mechanisms can lead to a red-shift or blue-shift in the PL peak energy depending on the materials involved. However, Olkhovets et al. (1998) reported that electron-phonon coupling contribution is significantly reduced in the QDs thus produce weaker dependence on temperature of the energy emission in the QD PL emission compared to bulk.

Figure 4(c) shows that the FWHM of the PL signal was also found to be temperature dependent and the linewidth broaden significantly with the increasing temperature for PbS QDs. Unlike isolated atoms, the carriers in the QDs are linked to the vibrational modes of the crystal lattice. The charges that localized in the QDs interact with lattice vibrations via phonons. At low temperature, PL spectrum reflects the distribution of the dot energies because of the charge carriers randomly frozen into the dot states. By increasing the temperature, the interaction between the electron and phonon leads to homogenous broadening of the optical linewidth (Turyanska et al. 2013). Another factor that also affect the PL linewidth is inhomogenous particle size distribution (Fomin et al. 1998). In previous studies of PbS QDs show thermally activated broadening of the linewidth at $T \sim 150$ K (Turyanska et al. 2007). This thermal activation is also observed for our PbS QDs sample in this research. In addition, the increasing of the FWHM with temperature is explained by the increase in exciton’s mobility with temperature (Gruning et al. 2004). As temperature increases, mobility of the excitons also increases, thus resulting in recombination (producing PL) with a broader (nonequilibrium) energy distribution.

Meanwhile, in Figure 4(d), the PL peak intensity shows that the PL intensity reaches maximum at the temperature of 60 K followed by monotonic decrease at $T > 80$ K. The increase of the PL intensity from 10 K to 60 K indicates that with the increasing T, carriers that are trapped in the defects of the QDs overcome the energy barriers and hence falling down to the QDs ground state. In other words, this is attributed to the thermal excitation of carriers out of QDs defect states followed by carrier relaxation into the ground state. In contrast, the PL intensity shows decreasing trend at the temperature $T > 80$ K. This quenching of the PL intensity at $T > 80$ K is associated with the excitation of carriers out of the QDs into non-radiative recombination centers. As temperature increases, the charge carriers gained more energy resulting in recombination of electrons and holes happened rapidly.
at a very high rate. This causing intensity to be diluted. Besides that, the quenching of PL emission intensity could partially due to the increasing number of surface defects caused by the ligand displacement (Turyanska et al. 2009).

POWER DEPENDENCE OF PL
Analyzing the dependence of integrated PL (IPL) as a function of excitation power density can provide information of the dominant carrier recombination mechanism (Hasbullah et al. 2009). In Figure 5, the integrated PL across the spectral region (~ 0.7-1.55 eV) is plotted on a log-log scale. Based on their gradients, $m$, the data in Figure 5 can be categorized into three groups, namely, $m \sim 1$, $1 < m < 2$, and $m \sim 2$. In this groups, they correspond to the condition where dominant carrier recombination mechanism is radiative ($m \sim 1$), mixture of radiative and nonradiative (associate with defects) ($1 < m < 2$), or nonradiative dominant ($m \sim 2$) (Hasbullah et al. 2009). In other study, Martini et al. (2011) reported that, in a double logarithmic scale, the integrated PL intensity increases linearly with the increasing excitation power density and can be described by the relation,

$$I_{PL} = \eta I_0^\alpha$$  \hspace{1cm} (2)  

where is the integrated PL intensity; and is the excitation power density (Jin et al. 1997). The exponent $\alpha$ depends on the radiative recombination mechanisms and $\eta$ is the coefficient related to the PL efficiency of the QDs and includes several effects as capture, absorption, ionization and recombination of excitons. As can be seen in Figure 5, the IPL increases linearly with the increasing power density. This similar behavior also can be seen in PbS QDs film and others semiconductor material (Marshall et al. 2016; Ru et al. 2003). However, our sample shows at all temperature measured, the gradient of the graphs is below one and closed to the unity where the radiative recombination is dominant. At 10 K, the slope is 0.69 and at this low temperature, the electron and hole recombine follows radiative recombination path due to excitonic recombination resulting from minimal thermal energy available for exciton dissociation. At 70 K and 150 K, $m \sim 0.86$ was observed at all excitation power density indicating that radiative recombination is dominant over the entire excitation range investigated. When we increase the temperature to 220 K and 300 K, the slope was increased to 0.91. Papagiorgis et al. (2016) also reported almost similar result where the slope for their PbS QDs solid is $m \sim 0.96$ above 200 K which exhibits a monomolecular behavior up to moderate excitation powers.

FIGURE 4. Temperature dependence of the PL properties for colloidal PbS QDs showing the PL peak energy, the FWHM and the PL intensity
Figure 6 shows the FWHM of the sample at different power excitation at low temperature (10 K) and at the room temperature (300 K). As can be seen in the Figure 6, the FWHM of the sample at 300 K are higher than at 10 K at all particular power excitation measured. As described previously, at higher temperature, the charge carriers become more mobile, thus broaden the PL linewidth with a nonequilibrium energy distribution. Furthermore, as the power excitation increases, the FWHM also increases for both temperatures (at 10 K and 300 K). This behavior is due to band filling effect where it happens by occupying the low energy states of energy bands by the particles.

CONCLUSION

In summary, colloidal PbS QDs capped with thiol ligands in aqueous solution was synthesized with stable photoluminescence in the near IR spectral range and the effect of temperature and power excitation towards the PL of PbS QDs was studied at the temperature range from 10 K to 300 K. The PL peak energy, PL peak intensity and FWHM of the sample are temperature dependent. Intensity of the PL increased at low temperature (10 K - 60 K). On the other hand, the intensity dropped when the temperature above 80 K due to activation of non-radiative recombination. We have observed a thermally activated increase the FWHM of our sample and explained this behavior in terms of the interaction between charge carriers and phonons at high temperature. We also found that the energy band gap of QDs is due to the size of QDs that related to quantum confinement effect. Recombination mechanism of charge carriers can be clearly observed by the measurement of power dependent towards PL emission.

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