First-Principles Study of Structural, Electronic and Thermoelectric Properties of Ni-Doped Bi\textsubscript{2}Se\textsubscript{3}  
(Kajian Prinsip Pertama tentang Sifat Struktur, Elektronik dan Termoelektrik bagi Bi\textsubscript{2}Se\textsubscript{3}, Ni-Terdop)

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ABSTRACT

Direct conversion of waste heat to electrical energy could address present energy challenges. Bi\textsubscript{2}Se\textsubscript{3} is one of few thermoelectric materials known to operate at room temperature. Comprehensive analysis using density functional theory was conducted to explore the effect of nickel doping on structural, electronic, and thermoelectric properties of Bi\textsubscript{2}Se\textsubscript{3}. Local density approximation (LDA) was used with an addition of spin-orbit coupling (SOC) and van der Waals interaction scheme consideration. Analysis of the effect of SOC was elaborated. It was found that nickel has changed the crystal structure of Bi\textsubscript{2}Se\textsubscript{3}. Nickel has also changed band structure and density of state that alter the thermoelectric performance. The decreased band gap has decreased the thermopower. However, it gives advantages to the improvement of electrical conductivity. Higher electrical conductivity has risen thermal conductivity. Despite the decreased thermopower and increased thermal conductivity, the higher electrical conductivity has improved the overall thermoelectric performance of Bi\textsubscript{2}Se\textsubscript{3} when nickel is introduced.

Keywords: Density functional theory; electrical conductivity; electronic properties; spin-orbit coupling; thermoelectricity

INTRODUCTION

Thermoelectric materials have received wide attention due to its ability to convert waste heat to electricity especially some of the materials that can effectively operating at room temperature (Bashir et al. 2018; Hasan et al. 2018). It could also potentially become one of the solutions for energy and environmental crisis like global warming (Saeed et al. 2014b). Its high-performance provides a great contribution to the society for its application in thermoelectric refrigeration (Mishra et al. 1997), radiation detector (Augustine et al. 2003) and environmental protector (Saeed et al. 2014b). The performance of thermoelectric material can be described by a dimensionless figure of merit (ZT):

$$ZT = \frac{S^2 \sigma T}{\kappa}$$

where $S$ is the Seebeck Effect or thermopower; $\sigma$ is electrical conductivity; $T$ is the temperature; and $\kappa$
is thermal conductivity. From the equation, a good thermoelectric material should possess high electrical conductivity and lower thermal conductivity. Bismuth selenide (Bi$_2$Se$_3$) is one of the excellent candidates for thermoelectric application due to the ability to operate at room temperature (Min et al. 2015), as compared to other thermoelectric materials that operate at elevated temperature (Chang et al. 2016; Tan et al. 2015, 2014). Additionally, Bi$_2$Se$_3$ is also a topological insulator, a new quantum state of matter that gives the thermoelectric material unique properties that could conduct electricity on its surface or edge, but still having insulating behaviour in the bulk (Xu et al. 2017). Among the topological insulators, Bi$_2$Se$_3$ is one of the most extensively studied because of its simplicity surface state which made up of a single Dirac cone at the Γ point (Aguilera et al. 2013). This discovery could make this material a promising candidate to be applied such as on spintronic devices (Mirihosseini & Henk 2012), topological quantum computing (Lin et al. 2013), and magnetic memories (Chis et al. 2012).

Despite the advantages of Bi$_2$Se$_3$, that can operate at room temperature, the thermoelectric performance at its pure form in room temperature is very low which has figure of merit around 0.01 - 0.05 (Kulsi et al. 2017; Min et al. 2015, 2013). Therefore, many attempts have been done to increase the thermoelectric performance of Bi$_2$Se$_3$. Many experimental methods have been conducted that show different performance which indicates that method gives influence on the performance. Scalable exfoliation strategy has improved Bi$_2$Se$_3$ performance up to 0.11 figure of merit (Sun et al. 2012). Pulsed magnetic field method increased the Bi$_2$Se$_3$ performance to 0.14 (Kulbachinskii et al. 2012). In addition, different types of nano-structure also influence the performance of Bi$_2$Se$_3$, such as nanoplates structure that gives Bi$_2$Se$_3$ figure of merit 0.14 (Ali et al. 2014). There are also many efforts to enhance the performance by using a doping method such as reported by Hor et al. (2009) that using calcium as a dopant and get figure of merit 0.07. Doped by using Cu and Fe increased the performance to 0.10 (Dun et al. 2015; Saeed et al. 2014b). There was also an effort to increase the power factor by doping with Ti and Cr (Cermak et al. 2018; Janiček et al. 2009).

This work will use another transition metal, Ni to increase the thermoelectric performance of Bi$_2$Se$_3$. Besides that, we intended to investigate the underlying mechanism of the enhanced thermoelectric performance by conducting a comprehensive analysis of the effect of Ni on the structural and electronic properties of Bi$_2$Se$_3$. We expected the role of Ni will reduce bond length and band gap that indicates increased electron participation.

**COMPUTATIONAL DETAILS**

The calculation was performed using the rhombohedral structure of Bi$_2$Se$_3$, with the space group . The initial structure of Bi$_2$Se$_3$ was taken from the experimental study by Nakajima (1963). Structural optimization was performed by using CASTEP within pseudo-potential with cut-off energy 400 eV and k-points 3×3×1 after convergence test was done. The optimization was done within local density approximation (LDA) exchange-correlation and Perdew-Burke-Ernzerhof revised for solids (PBEsol) approximation (Perdew & Zunger 1981). Van der Waals interaction scheme was considered. The structural properties that consist of bond length, lattice parameter and bond order were calculated by optimizing the structure.

Electronic properties were performed using WIEN2k (Blaha et al. 2001) computer code within full-potential linearized-augmented plane-wave (FP-LAPW) and exchange-correlation of LDA (Perdew & Zunger 1981). The electronic properties calculation which are band structure, the density of states (DOS) and charge density were performed with the inclusion of spin-orbit coupling (SOC) to analyse its effect on the material (Richards et al. 1981).

The calculation of thermoelectric properties was performed by Boltzmann Transport Properties (BoltzTraP) that implemented within the package of WIEN2K (Madsen & Singh 2006). The thermopower, electrical conductivity, thermal conductivity, and figure of merit (ZT) were evaluated over temperature from 300 to 500 K. The calculation was performed under the plane-wave cut off of 8.0 and k-points mesh of 1000.

**STRUCTURAL OPTIMIZATION**

In this work, the structural optimization was performed to stabilize the geometric cells and the structure of Bi$_2$Se$_3$ and Ni-doped Bi$_2$Se$_3$. Structural properties which are bond length, lattice parameter, and bond order were calculated. Figure 1 shows an optimization structure of Bi$_2$Se$_3$.

Bi$_2$Se$_3$ is a rhombohedral structure that consists of five-atom layers arranged along z-direction by the repeated atom sequences of Se(1) – Bi – Se(2) – Bi – Se(1). These layers are known as quintuple layer. The bonding within the quintuple layer is strong but much weaker between two quintuple layers that are bonded by van der Waals interaction. Table 1 shows the structural properties of Bi$_2$Se$_3$ and Ni-doped Bi$_2$Se$_3$.

Our result averagely consistent with experimental studies. Within LDA calculation, lattice a averagely overestimate only by 0.05%, while lattice c averagely
underestimates by 5.14%. Within GGA-PBESol calculation, lattice $a$ averagely underestimate by 0.19%, while lattice $c$ averagely underestimates by 0.57%. LDA calculation has more consistent value on lattice $a$, while GGA-PBESol has more consistent value with lattice $c$ when compared with experimental studies. Since the gap of consistency between LDA and PBESol is not having a large difference, LDA was chosen to optimize the structure of Ni-doped $\text{Bi}_2\text{Se}_3$ and to calculate the electronic and thermoelectric properties of undoped and doped $\text{Bi}_2\text{Se}_3$. LDA was chosen to save cost, time and energy since LDA has less parameter to calculate when compared with PBESol. Doping with Ni has changed the lattice parameter of $\text{Bi}_2\text{Se}_3$ with $a$ and $b$ decreasing by 2.44%, and $c$ decreasing by 10.28%. Nickel has made the whole structure of $\text{Bi}_2\text{Se}_3$ smaller by 14.31% from volume $404.269 \, \text{Å}^3$ to $346.399 \, \text{Å}^3$.

**FIGURE 1.** Optimized structure of $\text{Bi}_2\text{Se}_3$ and Ni-doped $\text{Bi}_2\text{Se}_3$

**TABLE 1.** Structural properties of $\text{Bi}_2\text{Se}_3$ and Ni-doped $\text{Bi}_2\text{Se}_3$

<table>
<thead>
<tr>
<th>Method</th>
<th>Reference</th>
<th>$a(\text{Å})$</th>
<th>$b(\text{Å})$</th>
<th>$c(\text{Å})$</th>
</tr>
</thead>
<tbody>
<tr>
<td>LDA</td>
<td>This work (Pure)</td>
<td>4.147</td>
<td>4.147</td>
<td>27.221</td>
</tr>
<tr>
<td>GGA-PBESol</td>
<td>This work (Pure)</td>
<td>4.137</td>
<td>4.137</td>
<td>28.493</td>
</tr>
<tr>
<td>LDA</td>
<td>This work (Ni-doped)</td>
<td>3.993</td>
<td>3.993</td>
<td>24.963</td>
</tr>
<tr>
<td>X-ray diffraction</td>
<td>Experimental (Boledzyuk et al. 2015)</td>
<td>4.147</td>
<td>4.147</td>
<td>28.681</td>
</tr>
<tr>
<td>X-ray diffraction</td>
<td>Experimental (Nakajima 1963)</td>
<td>4.143</td>
<td>4.143</td>
<td>28.636</td>
</tr>
</tbody>
</table>

The strength of the bond can be measured through the bond length for each atomic bonding in the structure. The bond length of $\text{Bi}_2\text{Se}_3$ and Ni-doped $\text{Bi}_2\text{Se}_3$ are tabulated in Tables 2 and 3, respectively.
The van der Waals length is 3.34 Å, which is the longest length shows the weakest bonding in the structure of Bi₂Se₃. The weak van der Waals bonding between adjacent quintuple layer has caused Bi₂Se₃ having topological surface states. The bond of Se(1) – Bi shows the shortest bond which has a length of 2.83 Å. Se (1) – Bi is the strongest bond in the structure which is located within the quintuple layer. Besides the contribution to the electron conductivity, the short bond length also contributes to the thermal conductivity due to its short-range potential. However, the significance of van der Waals in the structure is it can reduce phonon transmission throughout the material, thus, reduces thermal conductivity as reported by Park et al. (2015). The calculated bond lengths are relatively consistent with the experimental works. Se (1) – Bi underestimated by 4.95% and van der Waals length overestimated by 2.10%. The bond of Se (2) – Bi remains the same with the experimental works (Hagmann et al. 2017; Lind et al. 2005).

Nickel (Ni) atoms were doped at the van der Waals connecting atoms by substituting two atoms of Bi-lattice as shown in Figure 1. It shows the optimization structure of Ni-doped Bi₂Se₃. The presence of Ni atoms caused the bond length of Se (1) – Bi and Se (2) – Bi decreased. The van der Waals distance remains the same. The bond length of Ni-doped Bi₂Se₃ is tabulated in Table 3.

Table 3 shows that Se (1) – Ni is the strongest bond in the structure of Ni-doped Bi₂Se₃ due to its shortest bond length. Se (1) – Bi bond has decreased by 0.35% and Se (2) – Bi bond has decreased by 0.66%. A new bond has formed within the structure which is Se (1) – Ni bond and Se (2) – Ni. Decreased bond length has increased electron participation, thus increasing the electrical conductivity. Due to the unchanged distance of van der Waals interaction, the reduction of phonons transmission remained the same. However, due to the decreased bond length, the thermal conductivity would be increased.

The bond order is calculated to analyse the stability of the bond. The results of the calculated bond order of Bi₂Se₃ and Ni-doped Bi₂Se₃ are tabulated in Table 4.
Se (1) – Bi has decreased its ionic bonding by 14.91% which shows that nickel has weakened the bond. Bond of Se (2) – Bi remained as a single covalent bond. The new formed of the bond which are Se (1) – Ni and Se (2) – Ni have mix covalent-ionic type of bonding. Introducing nickel to the Bi-lattice, the ionic bonding has increased by 33.33% when compared between Se (1) – Ni and Se (1) – Bi. For Se (2) – X (X=Bi, Ni), the bond has changed from covalent bond to the mixed covalent-ionic bond.

### ELECTRONIC PROPERTIES

The band structure and density of states (DOS) were calculated by using LDA as the exchange-correlation with cut-off energy 8.0 and k-points 1000. The results were compared with the other works on both experimental and theoretical studies. The band gap of Bi$_2$Se$_3$ with the inclusion of SOC is 0.23 eV while without SOC is 0.01 eV as shown in Figure 2.

The band gap between without and with the inclusion of SOC shows a huge difference up to 82.35%. The SOC band gap is consistent with a small difference of 4.26% with experimental work by Martinez et al. (2017) (0.23 eV). However, the band gap underestimates with other works by difference around 18 - 31% (Han et al. 2014; Lawal & Shaari 2017; Valla et al. 2012; Wang & Zhang 2012). On the other hand, the band gap without SOC shows inconsistency between each other (Lawal & Shaari 2017; Wang & Zhang 2012.). This proved that

<table>
<thead>
<tr>
<th>Bond</th>
<th>Bond order</th>
</tr>
</thead>
<tbody>
<tr>
<td>Se (1) – Bi</td>
<td>1.000 0.570</td>
</tr>
<tr>
<td>Se (2) – Bi</td>
<td>1.000 1.000</td>
</tr>
<tr>
<td>Se (1) – Ni</td>
<td>1.000 0.760</td>
</tr>
<tr>
<td>Se (2) – Ni</td>
<td>1.000 0.760</td>
</tr>
<tr>
<td>Se (1) – Bi</td>
<td>1.000 0.485</td>
</tr>
<tr>
<td>Se (2) – Bi</td>
<td>1.000 1.000</td>
</tr>
</tbody>
</table>

Table 4 shows the bond order for each bond length of Bi$_2$Se$_3$. Se (1) – Bi have two bond orders which are 1.000 and 0.570 that indicates a mix of ionic-covalent type of bonding. Bond order of 1.000 indicates a single covalent bond, while 0.570 indicates ionic bond. Se (2) – Bi has a bond order of 1.000 that indicates a single covalent bond. 1.000 is the maximum bond order and the strongest bond of Bi$_2$Se$_3$ structure.

Nickel has made a significant change to the bond order of Bi$_2$Se$_3$. Most of the bond length has a mixed covalent-ionic type of bonding except for Se (2) – Bi. Bond

![FIGURE 2. a) Band structure of Bi$_2$Se$_3$ without SOC at energy range from 10 eV to -10 eV b) Band structure of Bi$_2$Se$_3$ with SOC at energy range from -10 eV to 10 eV c) Band structure of Bi$_2$Se$_3$ without SOC at energy range from 1.0 eV to -1.0 eV d) Band structure of Bi$_2$Se$_3$ with SOC at energy range from -1.0 eV to 1.0 eV](image-url)
SOC inclusion is necessary to calculate Bi$_2$Se$_3$, hence to manifest its true nature. Other theoretical works such as by Lawal and Shaari (2017) and Wang and Zhang (2012) show the same behaviour regarding this. Without the inclusion of SOC, Bi$_2$Se$_3$ could not manifest the band structure correctly compared to the inclusion of SOC.

Anti-crossing feature occurred around as a result of the effect of SOC on the band structure. This feature yields an inversion between conduction and valence band. As we can observe in Figure 2, with the inclusion of SOC the band structure formed a multi-valley band at maximum valence band instead of single-valley that shown by band structure without the inclusion of SOC. This inversion is known as spin-orbit splitting whereby the SOC splits the bands. The splitting energy occurs due to the Dresselhaus effect from the applied SOC. Without SOC, a gapless band structure would occur. SOC breaks the gapless band structure that resulting the Dresselhaus effect (Dresselhaus 1955). Besides that, the occurrence of a multi-valley band means there is degeneracy. This is the sign of high electrical conductivity.

Besides thermoelectric material, Bi$_2$Se$_3$ is also a topological insulator. One of the main properties of the topological insulator is the presence of the magnetic field at the absence of applied magnetic field from the macroscopic scale. This effect is known as the quantum Hall effect which is Hall effect that occurs at quantum state with the absence of the magnetic field (Novoselov et al. 2007). This effect gives unique properties to Bi$_2$Se$_3$ such as time-reversal symmetry protection and protection from non-magnetic impurities. Any non-magnetic impurities that put on the material would not change its properties due to its topological nature (Xu et al. 2017). The role of SOC is the correction method to calculate an additional magnetic field that occurs in the material. SOC gives spin-dependent force to the moving electron that yields the occurrence of the magnetic field. However, SOC could not produce the quantum Hall effect because it breaks time-reversal symmetry. Despite that, SOC induce a quantum spin Hall effect that can conserve the symmetry due to the reliance on the spin of the electron. Thus, instead of time-reversal symmetry, SOC provides protection based on the spin conservation symmetry. This is the main reason for the necessity of SOC in the calculation for Bi$_2$Se$_3$.

The effect of SOC is significant on the heavy element, which in this study, bismuth is one of the heavy elements. This is important for SOC to modify the electronic structure. In addition, a small band gap is another important factor for SOC to have a significant effect in the calculation (Parker & Singh 2011; Shi et al. 2015).

From the result of DOS without the inclusion of SOC as shown in Figure 3, it shows that Se-p has most population contribution at the valence band of electronic structure with the peak of 1.4 States/eV and energy range from 0 eV to -5.5 eV. Se-p partially hybridized with Bi-p and Bi-s. At the conduction band, Se-p has the most contribution with almost fully hybridized with Bi-p and density peak of 0.5 States/eV and energy range occupation from 0 eV to 5.0 eV.

![FIGURE 3. Density of states (DoS) of Bi$_2$Se$_3$ without SOC and inclusion of SOC](image-url)
From the DOS with the inclusion of SOC, the pattern of population contribution is almost the same but the density peak of Se-p is reduced to 1.0 States/eV at the valence band. Se-p possesses the most contribution at the conduction band with strong hybridization with Bi-p. The hybridization of Se-p and Bi-4p occupied at energy range from 0 to 5.5 eV. The occupation of electrons without SOC is higher which indicates higher electrical conductivity.

Band structure of Ni-doped Bi$_2$Se$_3$ is calculated as shown in Figure 4 to analyse the changing of band structure when Bi$_2$Se$_3$ is doped with nickel. Ni-doped Bi$_2$Se$_3$ has an indirect band gap of 0.025 eV. The band gap is reduced considerably by 89.58% from 0.24 eV. This finding is supported by Feng et al. (2017) and Saaed et al. (2014a). The increased degeneracy of hole at valence band can be observed at the range of -0.1 to 0.0 eV. This high degeneracy is due to the high level of doping which shows that nickel gives a significant effect on the band structure of Bi$_2$Se$_3$. High level of doping and degeneracy has shifted the nature of semiconductors to act more like semi-metal. Near the Fermi level, the band valley at the valence band has decreased. However, the degeneracy is increasing. The increased degeneracy indicates the increase of electrons carriers thus increasing the thermopower of Bi$_2$Se$_3$. High degeneracy and reduced band gap show the sign of higher electrical conductivity. High degeneracy shows there are multiple indirect band gap occurred, thus thermal conductivity is increased.

**FIGURE 4.** Band structure of Ni-doped Bi$_2$Se$_3$

**FIGURE 5.** Density of states of Ni-doped Bi$_2$Se$_3$
The density of states of Ni-doped Bi\textsubscript{2}Se\textsubscript{3} is calculated and the result is shown in Figure 5. Nickel gives considerable effect to the occupancy of holes at the valence band. Ni-f and Ni-s are most contributed at both valence and conduction band at the range of energy from 0 to -4 eV and 0 to 2 eV, respectively. Se-s and Se-d are strongly hybridized at valence and conduction band that indicates a strong covalent bond among Se atoms. The strong hybridization also occurred on Ni-s and Ni-f. The density of Ni-s and Ni-f reaches a peak of 6 states/eV. At the conduction band, the Ni-s and Ni-f density are far smaller than their respective density at the valence band. However, the occupation of electrons is higher than pure Bi\textsubscript{2}Se\textsubscript{3}. The density at the conduction band is increased by around 66.66% from ~0.6 to ~1.0. This shows the increase of electrical conductivity.

**THERMOELECTRIC PROPERTIES**

The thermoelectric properties of Bi\textsubscript{2}Se\textsubscript{3} and Ni-doped Bi\textsubscript{2}Se\textsubscript{3} are calculated and tabulated in Tables 5 and 6, respectively. The data from each table were used to plot a graph for each thermoelectric properties that consist of thermopower $S$, electrical conductivity $\sigma$, thermal conductivity $\kappa$ and figure of merit $ZT$ as shown in Figure 6. A comparison is made between Bi\textsubscript{2}Se\textsubscript{3} and its Ni-doped based on each correspondence thermoelectric properties.

**TABLE 5. Thermoelectric properties of Bi\textsubscript{2}Se\textsubscript{3}**

<table>
<thead>
<tr>
<th>T(K)</th>
<th>$S$ ($\mu V/K$)</th>
<th>$\sigma$ (S/cm)</th>
<th>$\kappa$ (W/mK)</th>
<th>$ZT$</th>
</tr>
</thead>
<tbody>
<tr>
<td>300</td>
<td>86.888</td>
<td>179.898</td>
<td>0.021</td>
<td>0.020</td>
</tr>
<tr>
<td>350</td>
<td>154.649</td>
<td>429.955</td>
<td>0.049</td>
<td>0.072</td>
</tr>
<tr>
<td>400</td>
<td>189.568</td>
<td>857.098</td>
<td>0.099</td>
<td>0.125</td>
</tr>
<tr>
<td>450</td>
<td>206.718</td>
<td>1501.186</td>
<td>0.173</td>
<td>0.167</td>
</tr>
<tr>
<td>500</td>
<td>214.300</td>
<td>2389.286</td>
<td>0.274</td>
<td>0.200</td>
</tr>
</tbody>
</table>

**TABLE 6. Thermoelectric properties of Ni-doped Bi\textsubscript{2}Se\textsubscript{3}**

<table>
<thead>
<tr>
<th>T(K)</th>
<th>$S$ ($\mu V/K$)</th>
<th>$\sigma$ (S/cm)</th>
<th>$\kappa$ (W/mK)</th>
<th>$ZT$</th>
</tr>
</thead>
<tbody>
<tr>
<td>300</td>
<td>-64.733</td>
<td>84993.700</td>
<td>0.611</td>
<td>0.170</td>
</tr>
<tr>
<td>350</td>
<td>-69.413</td>
<td>85160.014</td>
<td>0.744</td>
<td>0.193</td>
</tr>
<tr>
<td>400</td>
<td>-73.481</td>
<td>85283.742</td>
<td>0.868</td>
<td>0.212</td>
</tr>
<tr>
<td>450</td>
<td>-76.911</td>
<td>85498.860</td>
<td>0.999</td>
<td>0.228</td>
</tr>
<tr>
<td>500</td>
<td>-79.727</td>
<td>85825.755</td>
<td>1.135</td>
<td>0.239</td>
</tr>
</tbody>
</table>

The thermopower of Bi\textsubscript{2}Se\textsubscript{3} is increased from 300 to 500 K. Ni-doped Bi\textsubscript{2}Se\textsubscript{3} also shows the same trend. However, doped by nickel has decreased the thermopower of Bi\textsubscript{2}Se\textsubscript{3} over temperature. Positive values indicate that Bi\textsubscript{2}Se\textsubscript{3} possesses a hole as a carrier, while Ni-doped Bi\textsubscript{2}Se\textsubscript{3} has a negative value of thermopower which indicates electron as a carrier. Our thermoelectric results for pure Bi\textsubscript{2}Se\textsubscript{3} is in good agreement with the experiments by percentage difference around 4 - 26% over temperature (Kadel et al. 2011; Sun et al. 2015a, 2015b, 2012). At 300 K, the thermopower has decreased by 25.49% which is undesirable. This is due to the decreased band gap. The decreased band gap would decrease the thermopower (Feng et al. 2017; Lee & Mahanti 2012). In addition, the thermopower of Ni-doped is steadily increased as the temperature increased by a change rate of approximately 3% - 6% for every 50 K. The change has decreased when compared with pure Bi\textsubscript{2}Se\textsubscript{3} that has a change rate of around 3 - 78%.
conductivity. In our study, we calculated only electronic thermal conductivity. Electronic thermal conductivity is strongly influenced by electrical conductivity through Wiedemann-Franz law. The law states that increased electrical conductivity would also increase thermal conductivity (Bejan & Kraus 2003). Our results of thermal conductivity for pure Bi$_2$Se$_3$ are in good agreement with the experiments by percentage difference around 14 - 23% at temperature 350 and 400 K and underestimate around 48 - 66% at temperature 300, 350, and 500 K (Adam et al. 2018; Kang et al. 2017; Sun et al. 2015a). Doped with nickel has increased the thermal conductivity over temperature due to the increased electrical conductivity. Ni-doped has increased the thermal conductivity of Bi$_2$Se$_3$ from 0.021 to 0.611 W/mK at 300 K which is undesirable for thermoelectric material. However, the high electrical conductivity of Ni-doped Bi$_2$Se$_3$ has overcome the rising of thermal conductivity and thus improves the performance of Bi$_2$Se$_3$. Besides the influence from the electrical conductivity, the higher thermal conductivity also is due to the lattice contribution from indirect band gap at large maximum valence band to minimum conduction band at M to A. Without Ni-doped, the change rate of Bi$_2$Se$_3$ thermal conductivity is high up to 81% from 300 to 350 K. When the temperature increased to 400 K, the

The electrical conductivity can readily be derived from the transport distribution. Despite the disadvantage of the decreased band gap, it also gives advantages that contribute to the increased electrical conductivity. Decreased band gap would easily increase the electrons excitation from valence band to the conduction band, hence increase conductivity (Feng et al. 2017). Figure 5 shows electron occupancy is increased that indicates more electrons have excited and more electrical conductivity has occurred. Figure 6(b) shows increasing conductivity over temperature from 300 to 500 K for Bi$_2$Se$_3$, and its Ni-doped. Our results of electrical conductivity are in good agreement with the experiments by percentage difference around 1 - 15% over temperature (Kang et al. 2017; Saji et al. 2005; Sun et al. 2012a). Ni-doped has increased the electrical conductivity over temperature which is desirable for thermoelectric materials. Without Ni-doped, Bi$_2$Se$_3$ has an electrical conductivity change rate of around 66 - 82% from 300 to 400 K and decreased to 55% as the temperature reached 450 K. With Ni-doped, Bi$_2$Se$_3$ has a steady rate of change around 0.19 - 0.38%. Ni-doped has reduced the change rate of electrical conductivity of Bi$_2$Se$_3$ over temperature.

Thermal conductivity consists of two components which are electronic thermal conductivity and lattice thermal conductivity. In our study, we calculated only electronic thermal conductivity. Electronic thermal conductivity is strongly influenced by electrical conductivity through Wiedemann-Franz law. The law states that increased electrical conductivity would also increase thermal conductivity (Bejan & Kraus 2003). Our results of thermal conductivity for pure Bi$_2$Se$_3$ are in good agreement with the experiments by percentage difference around 14 - 23% at temperature 350 and 400 K and underestimate around 48 - 66% at temperature 300, 350, and 500 K (Adam et al. 2018; Kang et al. 2017; Sun et al. 2015a). Doped with nickel has increased the thermal conductivity over temperature due to the increased electrical conductivity.

Ni-doped has increased the thermal conductivity of Bi$_2$Se$_3$ from 0.021 to 0.611 W/mK at 300 K which is undesirable for thermoelectric material. However, the high electrical conductivity of Ni-doped Bi$_2$Se$_3$ has overcome the rising of thermal conductivity and thus improves the performance of Bi$_2$Se$_3$. Besides the influence from the electrical conductivity, the higher thermal conductivity also is due to the lattice contribution from indirect band gap at large maximum valence band to minimum conduction band at M to A. Without Ni-doped, the change rate of Bi$_2$Se$_3$ thermal conductivity is high up to 81% from 300 to 350 K. When the temperature increased to 400 K, the
thermal conductivity increased by 66%. The change rate decreasing down to 45% as the temperature reached 450 to 500 K. With Ni-doped, the change rate decreased around 13 - 20% over temperature from 300 to 500 K.

It can be observed in Figure 6(d) that the figure of merit (ZT) of Bi$_2$Se$_3$ is very low and not efficient to develop a thermoelectric device at room temperature. The results we obtained are in agreement with the experimental works by percentage difference around 2 - 36% (Kang et al. 2017; Kulsi et al. 2017; Min et al. 2015, 2013). Meanwhile, Ni-doped has improved the performance of Bi$_2$Se$_3$ by 750% for 300 K from 0.020 to 0.170. The ZT of Ni-doped Bi$_2$Se$_3$ increases steadily over temperature around with a change rate of around 5% - 12% for every 50 from 300 K.

CONCLUSION

The presence of nickel has changed the structure of Bi$_2$Se$_3$, forming a new bond and layer. The decreased bond length has increased electrical conductivity. The band gap of Bi$_2$Se$_3$ is in agreement with other available experimental studies with the inclusion of SOC. The SOC effect is necessary for the Bi$_2$Se$_3$ to show its band gap correctly due to its unique properties as a topological insulator. SOC effect give significant change such as multi-valley band, high degeneracy and multiple occurrences of indirect band gap. Bi$_2$Se$_3$ showed a poor thermoelectric material and thus, it make less efficient for application. However, Ni-doped has improved the performance of Bi$_2$Se$_3$ greatly and very suitable for room temperature thermoelectric devices. Ni has decreased the band gap Bi$_2$Se$_3$ and increase the occupation of electrons which indicates the increase in electrical conductivity and thermal conductivity. The decreased band gap has decreased thermopower. In addition, the occurrence of indirect band gap from a large maximum valence band to minimum conduction has also increased thermal conductivity. Despite the undesirable alteration of thermal conductivity and thermopower, the massive increase of electrical conductivity has overcome the deterioration.

ACKNOWLEDGEMENTS

This work was supported by the Ministry of Education (MOE) Malaysia under FRGS grant 600-IRM/FRGS 5/3 (031/2017) and Universiti Teknologi MARA (UiTM) and Universiti Teknologi Malaysia (UTM) for the facilities provided.

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Received: 10 April 2020
Accepted: 20 May 2020